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Systematics of alpha decay hindrance factors in doubly-even nuclei

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In present work, we have calculated the hindrance factors of 182 even-even alpha emitters using Preston's formulation of alpha decay probabilities and presented HFs systematics of 1⁻ and 3⁻ states in reflection asymmetric even-even quadruple-octupole deformed nuclides (A~216-230). The calculated HFs of both 1⁻ and 3⁻ states decrease with reduction in neutron number and this decrease is attributed to onset of intrinsic reflection asymmetry. There is a trend reversal for 1⁻ states at N=132 (218 Ra) and N=134 (220 Rn), which might be a possible indication of departure from static octupole deformation. Similarly, HFs systematics is discussed for $^{224-230}$ Th and $^{232-236}$ U isotopic chains along with 2⁺ states observed in daughter nuclei in N=132-146 isotonic chains.

Keywords: Even-even alpha emitters, Hindrance factors, Reflection asymmetry

1 Introduction

Alpha hindrance factor (HF), the ratio of experimental to theoretical partial half-lives of alpha transitions, is found to be an important tool in extracting nuclear structure information¹⁻⁴. Various theoretical techniques have been developed in order to understand the alpha-decay process and hence to calculate the penetrability of alpha particles through a barrier^{1,5}. The alpha transitions for which HF lies between 1 and 4, called favored transitions and take place between nuclear states having similar configurations and hence it is promising to ascertain both J^{π} and nucleonic configurations assignments for a given daughter (parent) state if those of parent (daughter) are known⁶. Similarly, the alpha HFs quantifies the correlation between nuclear wave functions of the initial state of parent and final state of daughter nuclei; larger wave function's overlap gives a lower HF³. The systematics of alpha-decay HFs can be used to deduce a variety of quantities like total alpha branching ratio, intensities of unobserved alpha groups and excitation energy of level in daughter nucleus⁶. In the present study, the spin-independent part of Preston's equations¹ have been used for the calculations of alpha decay probabilities. This formalism contains radius of the daughter nuclide, r_0 , as a free parameter. By setting HF=1 for the groundstate to ground-state alpha branch for an even-even nuclide², the radius parameter for the daughter nuclide can be deduced⁷ that can be used to deduce alpha hindrance factors for alpha-fed excited states in eveneven nuclides. We have calculated HFs of 182 eveneven alpha emitters in the framework of Preston's formulation¹ by using ALPHAD_RadD program⁸ and present the systematics of alpha HFs with daughter neutron number in reflection asymmetric (RA) mass region (A=216-230) i.e. for quadruple-octupole deformed nuclei. Additionally, the systematics of HFs for 2⁺ states observed in N=132-146 isotonic chains is also presented

2 Methodology

In the present study, a well established Preston's spin-independent formulation¹ with only essential steps described here, has been used. In this formulation, a preformed alpha particle is considered to be moving inside a nucleus having rectangular potential field of depth of V_0 ; $V_0 = \text{constant}$ for $r < r_0$, where r is the distance from the center of the product nucleus and r_0 is the radius of the product nucleus. The field beyond effective nuclear radius was assumed to be generated by a Coulomb potential $(2Ze^2/r, \text{ where } Z \text{ is the atomic number of product nucleus and e is the elementary charge) between alpha particle and daughter nucleus. The solution of the time dependent Schrodinger equation is assumed to have a form¹:$

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 $u = \psi(x, y, z) \exp(-iEt / \hbar)$

The wave-function ψ should obey following boundary conditions (a) at r = 0 at ψ is finite, (b) at $r = r_0$, ψ and $d\psi/dr$ are continuous, (c) ψ represents as outgoing wave for $r > r_0$. The Schrödinger equation for alpha particle in afore said potential field can be written as:

$$\frac{d^2 X_l}{dr^2} + \frac{2m}{\hbar^2} \left[E - V - \frac{\hbar^2 l \left(l + 1 \right)}{2mr^2} \right] X_l = 0$$

where, $E = E_{\alpha} - \frac{1}{2}i\hbar\lambda$; E_{α} is the energy of α -particle

and λ the time constant; are complex eigen values. In the interior of the nucleus with V=U, we have

$$X_{l}^{i} = \left\{\frac{2m(E-U)}{\hbar^{2}}\right\}^{\frac{1}{4}} r^{1/2} J_{l+\frac{1}{2}} \left[\left\{\frac{2m(E-U)}{\hbar^{2}}\right\}^{\frac{1}{2}} r\right]$$

where, J denotes Bessel functions and superscript i refer to interior of the nucleus. The solution X_i^i represents a standing wave and imaginary part of E related to the leak of alpha particle through the potential barrier.

3 Input Parameters

In the present work, the daughter radius parameter (r_0) is the main input used to calculate the HFs of alpha-fed excited states in even-even nuclides. In order to deduce r_0 a set of experimental quantities

such as Q_{α} value, parent nuclide's half-life (T_{1/2}), total alpha–decay branching ratios (% α), and alpha intensities (I_a) are used. The Q_{α} energies are taken from recent atomic mass evaluation of M. Wang et al.⁹, total alpha-decay branching ratios and half-lives of parent nuclides are taken from the ENSDF database¹⁰ supplemented by recent data from literature.

In order to calculate HFs of various excited states, the recently developed ENSDF analysis code namely, ALPHAD_RadD⁸, which is based on Preston's equations for alpha decay transition probabilities¹, has been used. This program can also be used to deduce HFs in odd-A and odd-odd nuclei by employing recently evaluated⁷ list of even-even daughter radius parameters.

4 Results and Discussion

In Reflection Asymmetric even-even quadrupleoctupole deformed nuclides (A~216-230), two separate bands with opposite parity i.e. $I^{\pi}=0^+$, 2^+ , 4^+ ... and $I^{\pi}=1^-$, 3^- , 5^- ... are generally observed⁴. In present paper, we studied the systematics of alpha HFs for 1⁻ and 3⁻ states with daughter neutron number in above said RA mass region (A~216-230). The results of HF systematics of 1⁻ and 3⁻ states observed in ²¹⁸⁻²²⁶Ra isotopic chains with daughter neutron number are presented in Fig.1(a). From this figure, it is clear that, the HFs of both 1⁻ and 3⁻ states smoothly decreasing with reduction in neutron number. This decrease of HFs is attributed to onset of



Fig. 1(a-d) — HF's systematics with daughter neutron number for ²¹⁸⁻²²⁶Ra, ²¹⁸⁻²²²Rn, ²²⁴⁻²³⁰Th and ²³²⁻²³⁶U nuclides.



Fig. 2(a-i) — HF's systematics of 2⁺ states observed in daughter nuclei in N=132-146 isotonic chains.

static Quadruple-Octupole deformation, but there is a trend reversal at N=132 (218 Ra), which might be a possible indication of departure from static Quadruple-Octupole deformation³.

Similar trend reversal in HFs systematics of 1⁻ state at N = 134 is also observed for ²¹⁸⁻²²²Rn nuclides (Fig.1 (b)), but such reversal could not be observed in ²²⁴⁻²³⁰Th (Fig.1(c)) as experimental data for 1⁻ state in lower mass region is not accessible. The HF systematics of 1⁻ state in ²³²⁻²³⁶U nuclides (Fig.1(d)) shows a minimum at N=142 and may be due to similar shape transition as discussed for Ra and Rn isotopic chains.

Additionally, we have also presented the systematics of HFs of 2^+ states observed in N=132-146 isotonic chains as shown in Fig. 2(a-i). From Fig. 2(a-c), a smooth decrease in HFs of 2^+ states with increase in proton number is observed till Z=88 in N=130, 132 and till Z=90 in N=134 isotopic chain. This smooth decrease in HFs indicates more probable alpha penetration through coulomb barrier and hence could be attributed to decrease in stability beyond Z=82 shell closure. The abrupt increase of HF at

Z=90 in N=130, 132 and at Z=92 in N=134 isotopic chain is still an open problem. Although there is increase in HFs in other isotonic chains (N=136-146) as shown in Fig. 2(d-i), but minima corresponding to certain Z number could not be identified due to inaccessibility of experimental data for these isotonic chains. On the basis of above systematics (Fig. 2(a-i)), we suggest that all the experimentally observed alpha decays branches (0⁺ to 2⁺) are favored.

5 Conclusions

The HFs systematics of 1⁻ and 3⁻ states observed in RA mass region is presented. The HFs of both 1⁻ and 3⁻ states decrease with reduction in neutron number and this decrease of HFs could be attributed to onset of intrinsic reflection asymmetry, but there is a trend reversal for 1⁻ state at N=132 (218 Ra) and N=134 (220 Rn), which might be a possible indication of departure from static octupole deformation. A smooth decrease in HFs of 2⁺ states with proton number is observed N=130, 132 and N=134 isotonic chains, which indicates the enhanced alpha penetration probability beyond Z=82 shell closure.

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